

Numerical Studies of Collapses in Turbulence in Nonlinear Schrödinger Equation

N. Vladimirova and P. Lushnikov

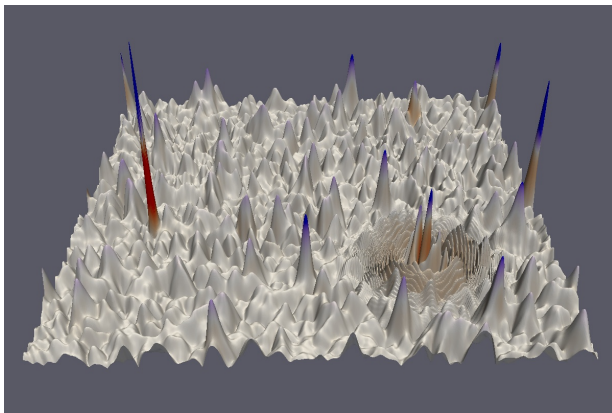
University of New Mexico

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Nonlinear Schrödinger equation (NLS)

Collapses in
NLS Turbulence

N. Vladimirova and
P. Lushnikov



$$i\psi_t + \nabla^2\psi + |\psi|^2\psi = 0$$

$$i\psi_t + (1 - i\epsilon\alpha)\nabla^2\psi + (1 + i\epsilon\beta)|\psi|^2\psi = i\epsilon\beta\psi$$

application: propagation of light through nonlinear media;
 $\psi(x, y, t)$ — envelope of electric field

NLS

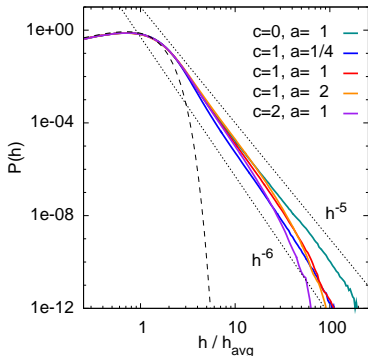
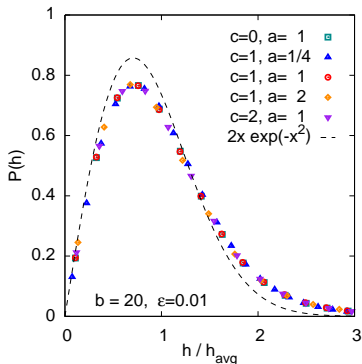
Turbulence

Collapses

Extra

Distribution of $|\psi|$ in the field

Notations: $h \equiv |\psi|$, $h_{avg} = \langle |\psi| \rangle \propto \sqrt{N}$, $N = \int |\psi|^2 d^2r$.



Turbulent background:

- ▶ pdf scales with h_{avg}
- ▶ universal shape for $h \sim h_{avg}$

Tails:

- ▶ power-law (?) for $h \gg h_{avg}$
- ▶ depend on a, c , not on b

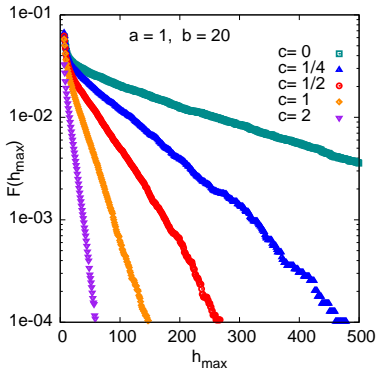
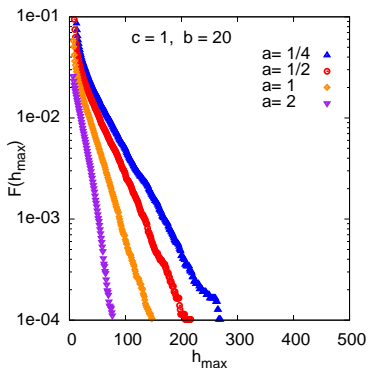
Other field quantities scale with h_{avg}

- Spatial correlation length ($\sim h_{avg}$) ←
- Correlation time ($\sim h_{avg}^2$) ← depend on a, c , not on b
- Spectra of $|\psi|_k^2$ ($\sim h_{avg}$) ←

When rescaled to h_{avg} ,

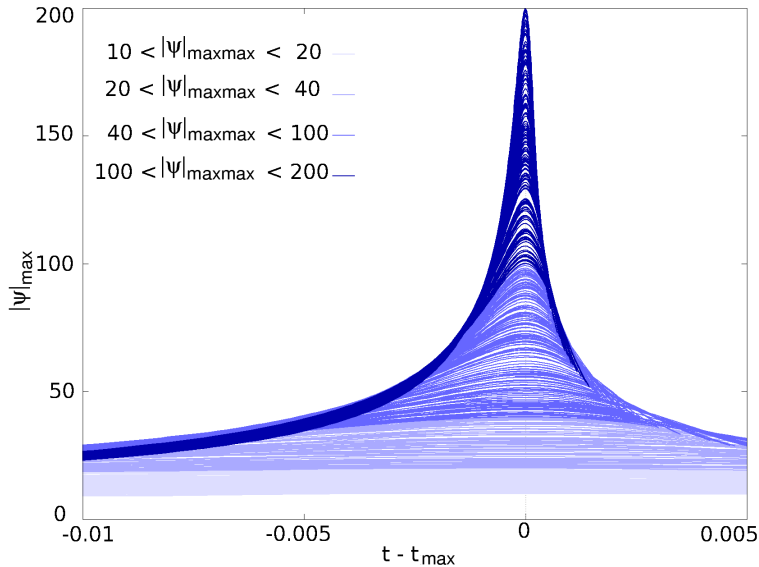
The shape of spatial correlation function is universal.

The shapes of temporal correlation function and spectra
(slightly) depend on a, c .

Frequency of collapses with $h > h_{\max}$.

- ▶ Frequency of collapses with $h > h_{\max}$ per unit area: $F(h_{\max})$.
- ▶ The dependence is exponential.
- ▶ The exponent depends on a , c , and b .
- ▶ No obvious scaling with h_{avg} .

Universality of collapses in turbulence: $a=1, c=0$



Collapses in
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NLS

Turbulence

pdf of $|\psi|$

... side note

pdf of collapses

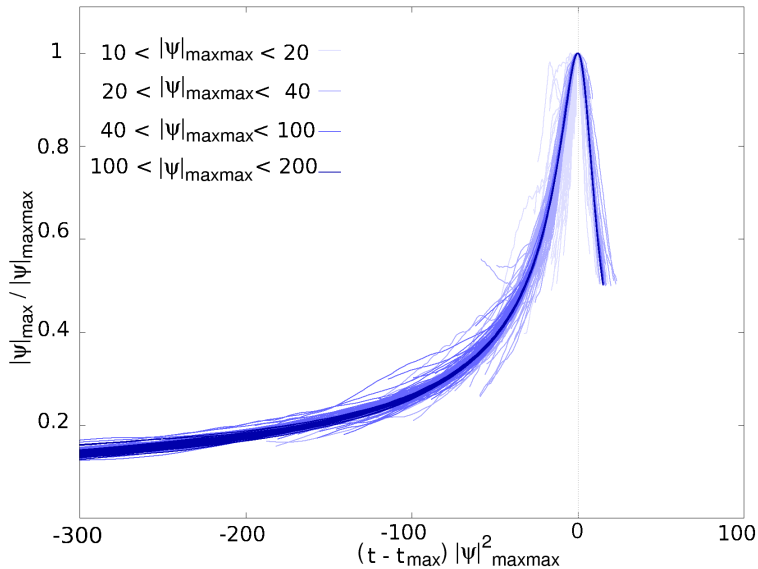
collapse similarity

$F(h_{\max}) \rightarrow pdf(\psi)$

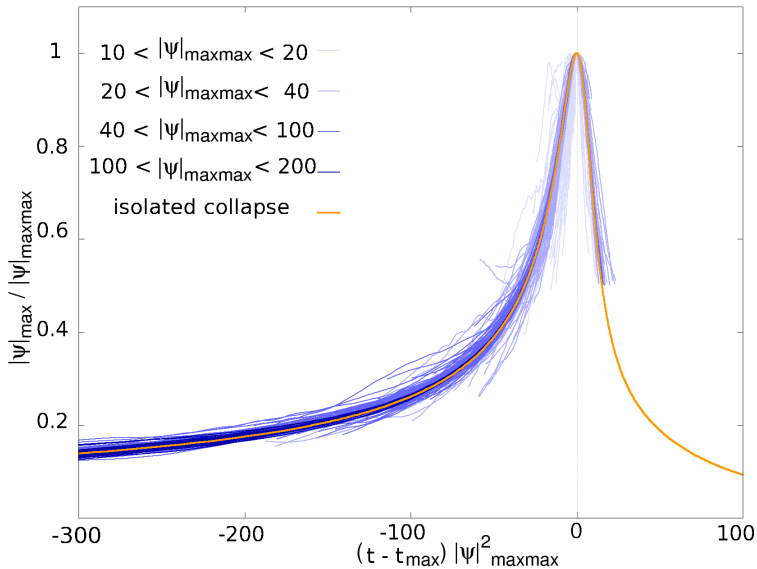
Collapses

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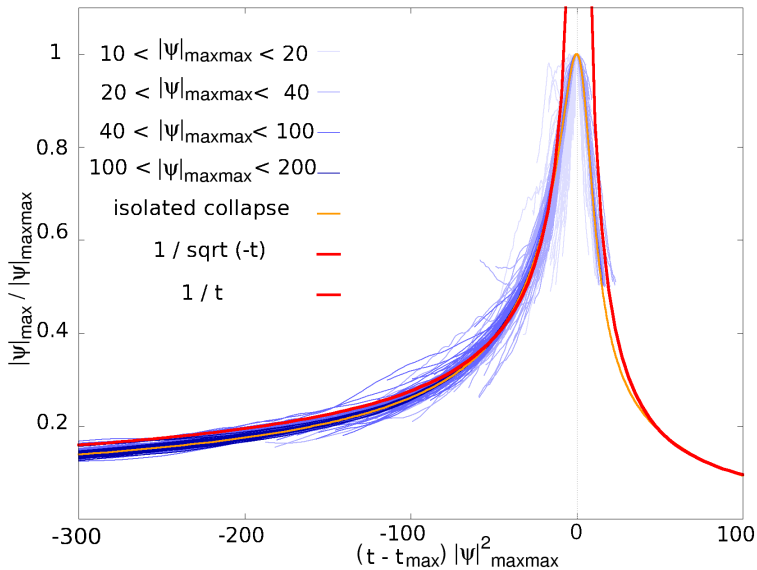
Universality of collapses in turbulence: $a=1, c=0$



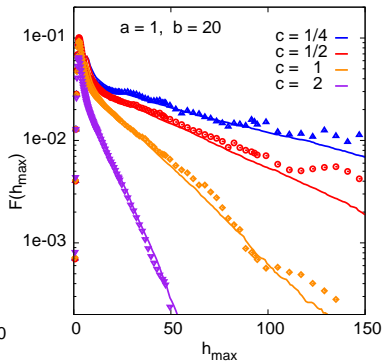
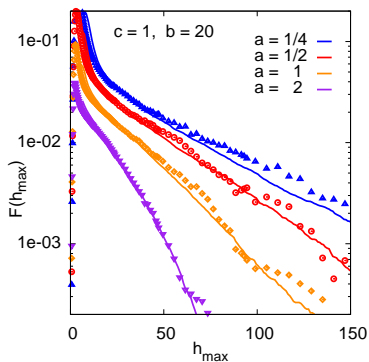
Universality of collapses in turbulence: $a=1, c=0$



Universality of collapses in turbulence: $a=1, c=0$



Connecting PDF of $|\psi|$ to frequency of collapses



- ▶ assume similarity among collapses, $h \sim (t_c - t)^{-\frac{1}{2}}$
- ▶ assume known frequency of collapses above h_{\max} : $F_{\max}(h_{\max})$
- ▶ conclude that PDF in the field $P(h) \sim h^{-5} F_{\max}(h)$
- ▶ Solid lines: $F(h_{\max})$. Points: $0.012 h^5 P(h)$.

NLS

Turbulence

pdf of $|\psi|$

... side note

pdf of collapses

collapse similarity

$F(h_{\max}) \rightarrow pdf(\psi)$

Collapses

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Diagnostics of collapses

- ▶ In slow variables $\rho = \frac{r}{L}$ and $\tau = \int L^{-2}(t')dt'$

$$\psi(r, t) = \frac{1}{L} V(\rho, \tau) e^{i\tau + i\gamma(\tau)\rho^2/4} \quad (1)$$

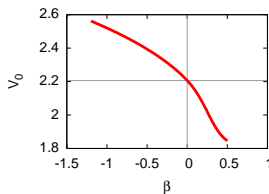
critical NLS becomes

$$\nabla^2 V + V^3 - V + \frac{1}{4}\beta\rho^2 V = 0. \quad (2)$$

- ▶ Solve (2) for each β , tabulate $V_0(\beta)$ at the center.
- ▶ Taylor expansion of (2) at the center and (1) give

$$L = L(|\psi|(0), |\psi|''_{rr}(0)) \Rightarrow V_0 \Rightarrow \beta; \quad \gamma = 2L^2\phi''_{rr}(0)$$

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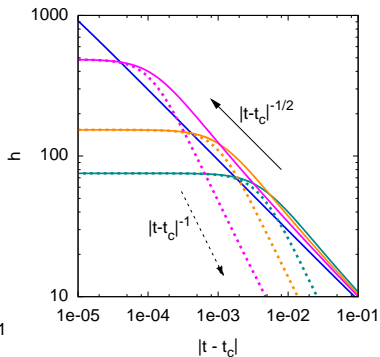
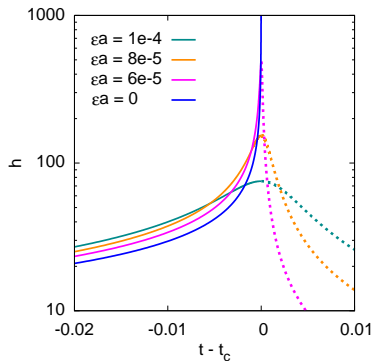
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Collapse growth and decay

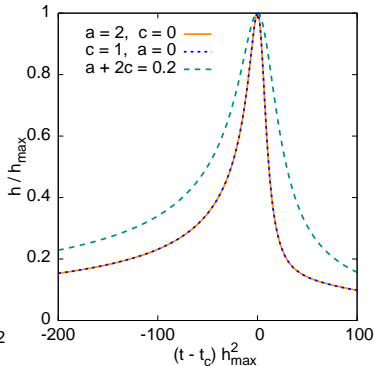
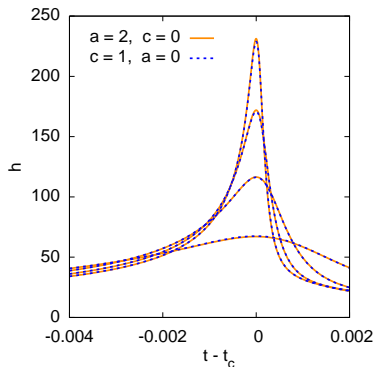


Evolution of collapses with $h_0 = 2.8$, $r_0 = 1$ with $c = 0$, $b = 0$.

Growth rate depends on a and c , decay rate does not.

Notation: $h \equiv |\psi|_{r=0}$.

Rescaled evolution: $(a + 2c)$ similarity



Left:

NLS coordinates

collapses with $\epsilon = 0.01$, $a + 2c = 2$, and different h_0 .

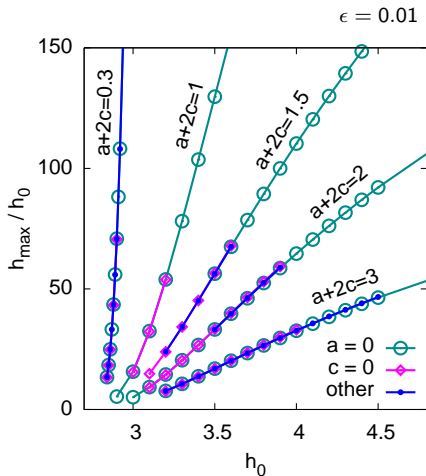
Right:

coordinated rescaled with h_{\max} ;

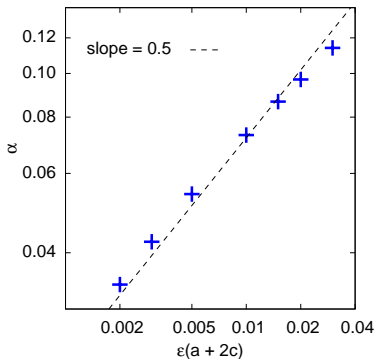
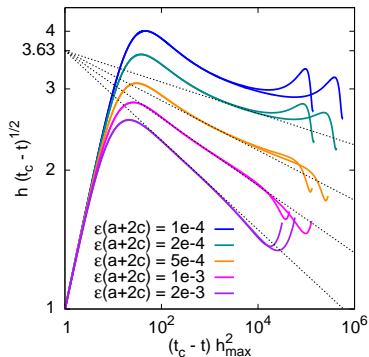
same curve for collapses with $a + 2c = 2$,

different for $a + 2c = 0.2$

Maximum height of collapses: $(a + 2c)$ similarity



Growth rate of collapses: $(a + 2c)$ similarity



$$h \propto (t_c - t)^{-(\frac{1}{2} + \alpha)}, \quad \alpha = \alpha(a + 2c)$$

Left: $h(t)$, compensated by $(t_c - t)^{-\frac{1}{2}}$.

Right: coefficient α determined from dotted lines on the left panel.

NLS

Turbulence

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growth and decay
rescaled $h(t)$
maximum height
growth rate
ODE model
ODE terms
...side note
ODE test

Extra

ODE model for collapse growth (Fibich and Levy, 1998)

$$L^2 \beta_t = -k_1 - k_2 \beta + k_3 (LL_t)^2 - \nu(\beta)$$

$$L^3 L_{tt} = -\beta$$

$$\nu(\beta) = 45.1 e^{-\pi/\sqrt{|\beta|}}$$

$$k_1 = 2\epsilon(a + 2c) \frac{N}{M}$$

$$k_2 = \epsilon a \frac{P - N}{M}$$

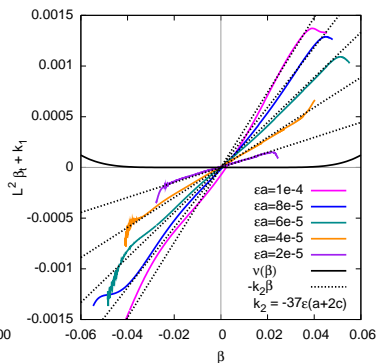
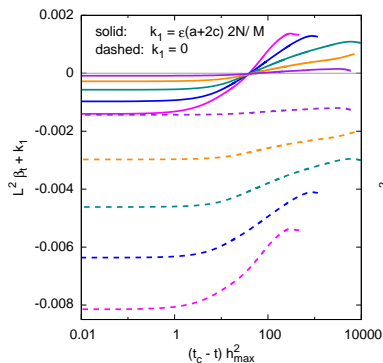
$$k_3 = \epsilon a \frac{P - N - 2M}{M}$$

$$N = 2\pi \int \rho R^2 d\rho \approx 11.69$$

$$P = 2\pi \int \rho^3 R_\rho^2 d\rho \approx 18.65$$

$$M = \frac{\pi}{2} \int \rho^3 R^2 d\rho \approx 3.46$$

Test of ODE model: term-by-term comparison



$$L^2 \beta_t = -k_1 - k_2 \beta + k_3 (L L_t)^2 - \nu(\beta)$$

$$L^3 L_{tt} = -\beta$$

Data suggest: k_1 term describes first order effects well;
 $k_2 \approx -37\epsilon(a+2c)$ instead of original $k_2 \approx 2\epsilon a$.

NLS

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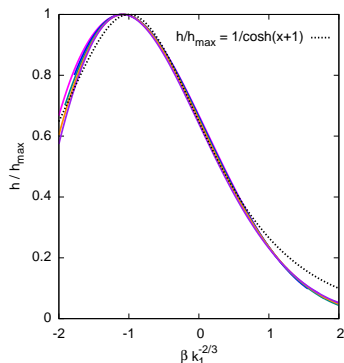
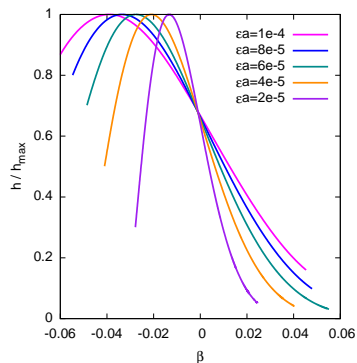
ODE terms

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Side note on ODE model



In slow variables keep first two terms of ODE model:

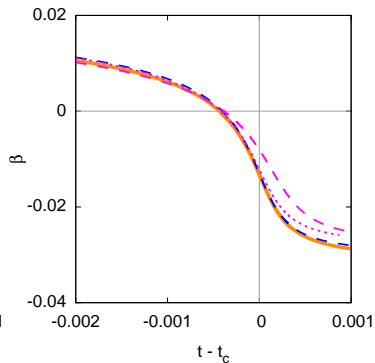
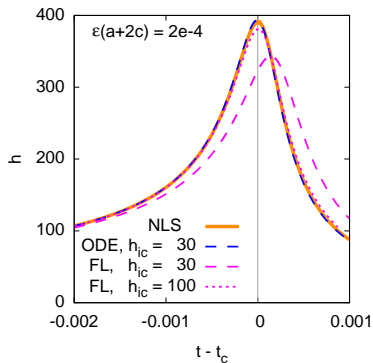
$$\beta_{\tau} = -k_1 - k_2\beta$$

$$L_{\tau\tau} - 2L_{\tau}^2/L = -\beta L$$

Results in $\frac{h}{h_{\max}} = \frac{L_{\max}}{L} \approx \cosh^{-1} \left[- \left(\tilde{\beta} - \tilde{\beta}_{\max} \right) |\tilde{\beta}_{\max}|^{\frac{1}{2}} \right],$

where $\tilde{\beta} = \beta k_1^{-\frac{2}{3}}$ and $\tilde{\beta}_{\max} = -0.92$

Test of ODE model: collapse evolution



NLS

Turbulence

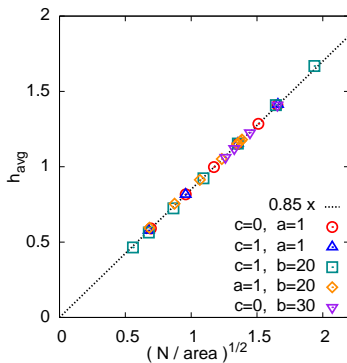
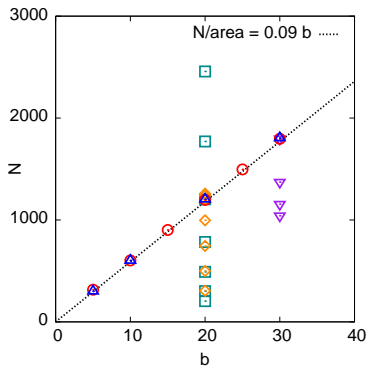
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Parameter space: a , c , and b ($\epsilon = 0.01$).



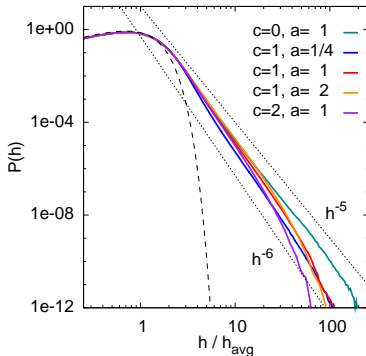
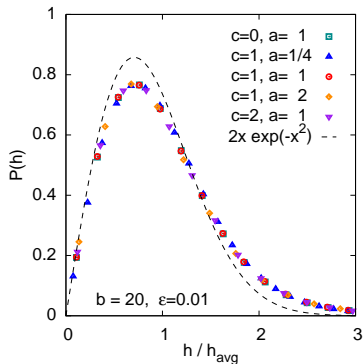
▶ Notations: $h \equiv |\psi|$, $h_{\text{avg}} = \langle |\psi| \rangle$, $N = \int |\psi|^2 d^2r$.

▶ No obvious scalings with a , c , or b ...
 ...except $N \propto b$ and $h_{\text{avg}} \propto \sqrt{N}$.

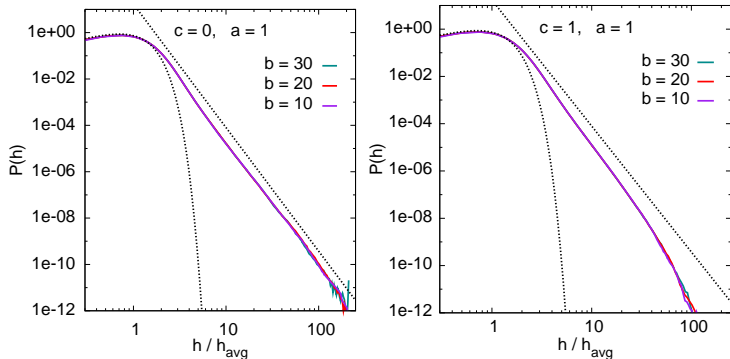
▶ We will show that h_{avg} well describes background turbulence.

▶ What determines h_{avg} ?

Distribution of $|\psi|$ in the field scales with h_{avg}

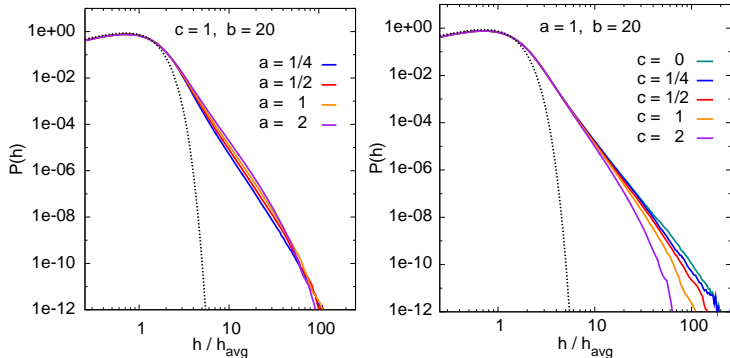


- ▶ PDF of $|\psi|$ scales with h_{avg} .
- ▶ Universal shape for $h \sim h_{avg}$, power-law (?) tails for $h \gg h_{avg}$.
- ▶ The tails do not depend on b .
- ▶ The tails do depend on a and c .

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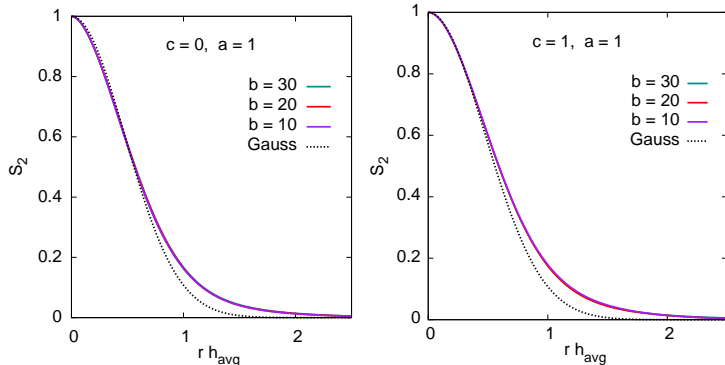
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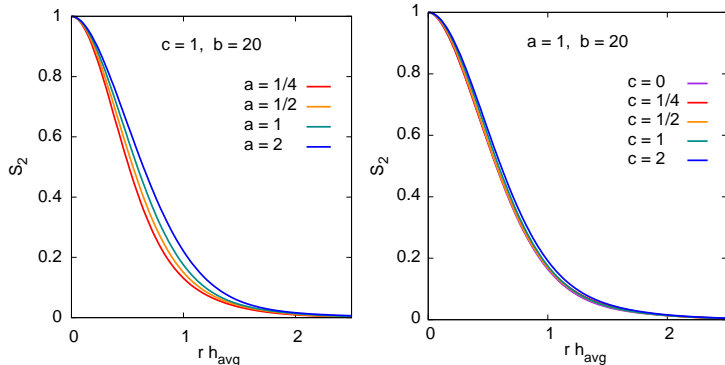
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Spatial correlation function scales with h_{avg}



- ▶ Spatial correlation function scales with h_{avg} .
- ▶ Correlation length does not depend on b .
- ▶ Correlation length depend on a and c .
- ▶ The shape, when rescaled to correlation length, is universal.

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NLS

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parameter space

pdf of psi

spatial corr fun

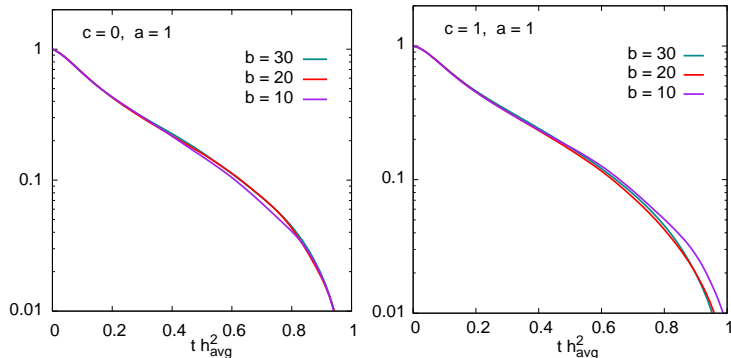
time corr fun

spectra

critical collapse

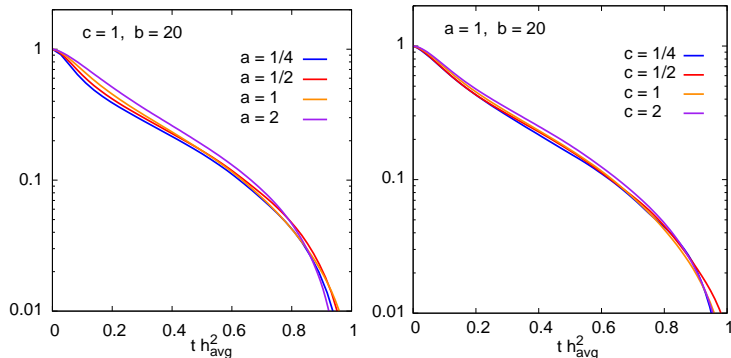
beta and gamma

Temporal correlation function scales with h_{avg}^2



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- ▶ Correlation time does not depend on b .
- ▶ Correlation time depend on a and c .

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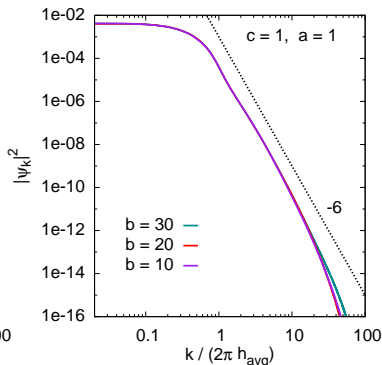
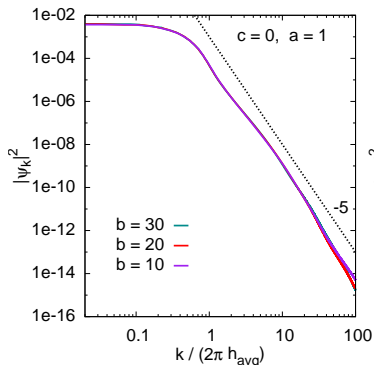
spatial corr fun

time corr fun

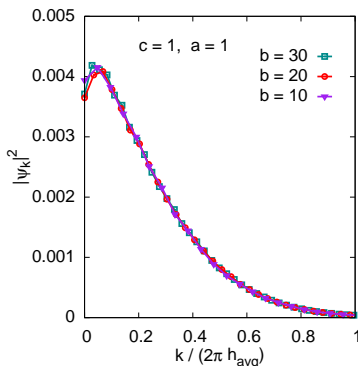
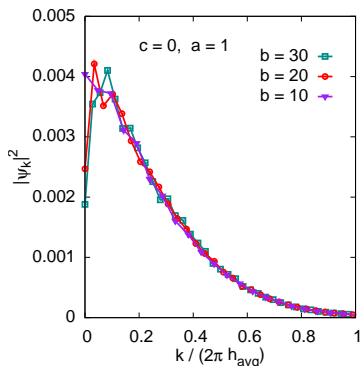
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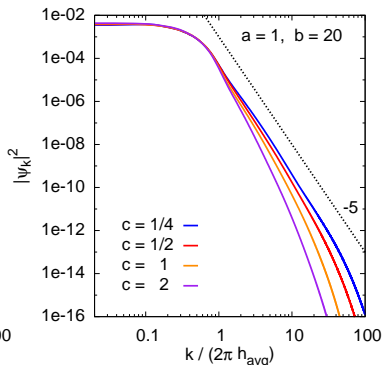
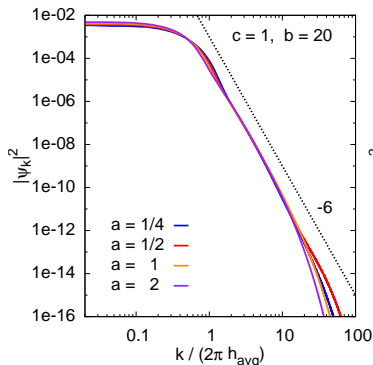
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Energy spectra scale with h_{avg} 

- ▶ Spectra of $|\psi|^2_k$ scale with h_{avg} .
- ▶ The tails do not depend on b .
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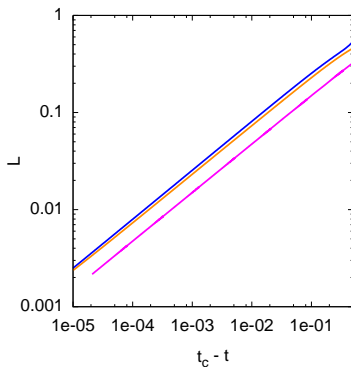
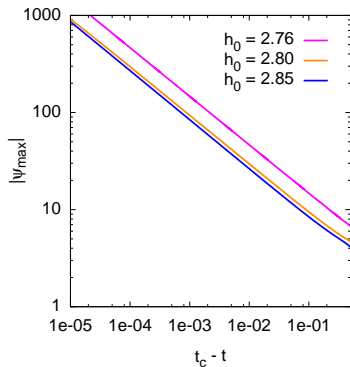
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Test: critical collapses



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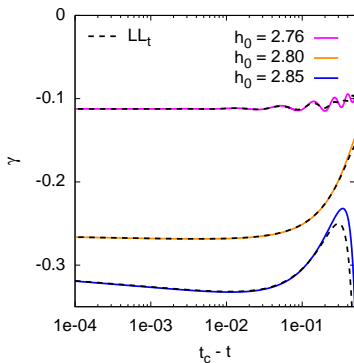
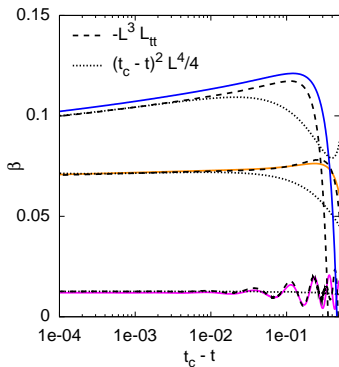
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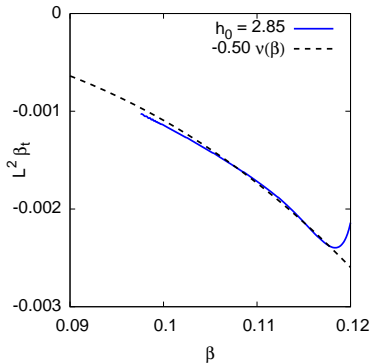
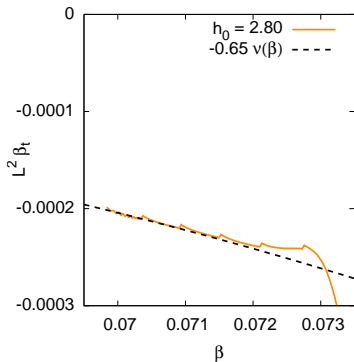
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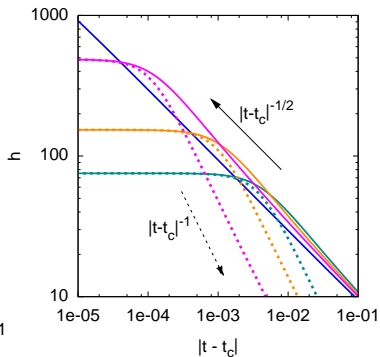
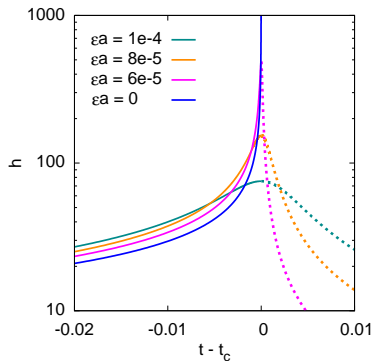
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Beta and gamma with stabilization



Evolution of collapses with $h_0 = 2.8$, $r_0 = 1$ with $c = 0$, $b = 0$.

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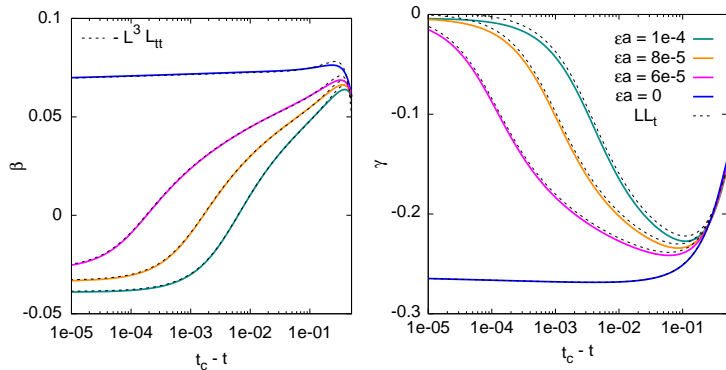
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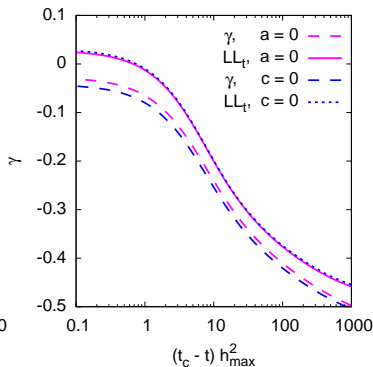
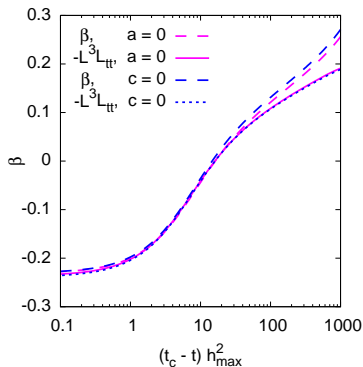
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Beta and gamma with stabilization



Here, $\epsilon(a + 2c) = 0.01$. Notice that $(a+2c)$ similarity is more pronounced in L and its derivatives rather than in β and γ .